# Introduction to microwave modeling Motivations for SMRT



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#### Context

Many snow microwave emission models have been developed over the last 30-40 years. The most "generic" models are widely used by the PM community: HUT, MEMLS, DMRT-QMS, DMRT-ML, ...

A few snow radar backscatter models have been developed by the AM community (side looking radar and nadir altimetry).

- several in specific studies
- DMRT-QMS (L. Tsang's group) is dual mode
- MEMLS has been extended to active mode in mid 2010's

In this introduction lecture:

- Why such a diversity?
- Is this diversity apparent or profound?
- Is this diversity beneficial or counter-productive for the community?
- What about the dual mode ? Good or bad ?
- Why a new model?

# Microwave model ingredients

- Snowpack
- Ground / sea ice / lake ice
- Forest, Atmosphere, ...



- Thermal Emission
- Scattering and absorption processes in the volumes
- Reflection and refraction at the surface/interfaces
- Inter-layer interferences

(e.g. ice crust, L-band, ...)



Snow is a **dense media** from the perspective of EM waves:

- Scattered by many particles  $\rightarrow$  change effective incident field
- Multiple scattering between particles
- → Concept of effective permittivity and Born approximation(s)





- Multi-species (e.g. wet snow)

In EM, sparse is up to 1 % frac vol, dense is >1%.

### Microwave model ingredients

Models differ in the ingredients and how detailed is each component described e.g. HUT (Snow + atmosphere) versus DMRT-ML (snow) + RTTOV (atmosphere)

Other constraints:

- frequency range

For typical snowpack:

No scattering	1st order scattering	Multiple « Rayleigh » Scattering	Multiple « Mie » Scattering	
	 1 GHz	 10 GHz	 100 GHz 2	200 GHz
Wave theory Transfer		Radiative Transfer		

SMRT is definitely a RT model. The following is about RT models

Other constraints (cont.):

- Application context  $\rightarrow$  performance, adjoint needed, ...
- language
- license
- ecosystem around the model, documentation, support, training, ...
- collaboration network, institutional constraints, community

There are many good reasons for different models.

But: our community is not so 'big'. Question: are the differences profound or superficial ?

The following is mostly based on « Are existing snow microwave emission models so different ? », Picard et al. AGU 2015

# Microwave model ingredients

Radiative transfer models in general:

The radiative transfer equation

$$\mu \frac{\partial \mathbf{I}(\mu,\phi,z)}{\partial z} = -\boldsymbol{\kappa}_{\mathrm{e}}(\mu,\phi,z) \mathbf{I}(\mu,\phi,z) + \frac{1}{4\pi} \iint_{4\pi} \mathbf{P}(\mu,\phi;\mu',\phi',z) \mathbf{I}(\mu',\phi',z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \mathbf{I}(\mu',\phi',z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \mathbf{I}(\mu',\phi',z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \mathbf{I}(\mu,\phi,z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \, \mathbf{I}(\mu,\phi,z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \, \mathbf{I}(\mu,\phi,z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \, \mathbf{I}(\mu,\phi,z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \, \mathbf{I}(\mu,\phi,z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \, \mathbf{I}(\mu,\phi,z) \, d\Omega' + \boldsymbol{\kappa}_{\mathrm{a}}(\mu,\phi,z) \, \alpha T(z) \, \mathbf{I}(\mu,\phi,z) \, \mathbf{I}(\mu$$

accompanied with boundary conditions:

$$\mathbf{I}^{(l)}(\mu < 0, \phi, z_{l-1}) = \mathbf{R}^{\text{spec,top},(l)}(\mu) \mathbf{I}^{(l)}(-\mu, \phi, z_{l-1}) + \frac{1}{2\pi} \iint_{2\pi,\mu' > 0} \mathbf{R}^{\text{diff,top},(l)}(\mu, \mu', \phi - \phi') \mathbf{I}^{(l)}(\mu', \phi', z_{l-1}) d\Omega' + \mathbf{T}^{\text{spec,bottom},(l-1)}(\mu) \mathbf{I}^{(l-1)}(\mu, \phi, z_{l-1}) + \frac{1}{2\pi} \iint_{2\pi,\mu' < 0} \mathbf{T}^{\text{diff,bottom},(l-1)}(\mu, \mu', \phi - \phi') \mathbf{I}^{(l-1)}(\mu', \phi', z_{l-1}) d\Omega'$$

### **Microwave model ingredients**

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The radiative transfer equation

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Accompagnied with boundary conditions:

$$\mathbf{I}^{(l)}(\mu < 0, \phi, z_{l-1}) = \mathbf{R}^{\text{spec}, \text{top}, (l)}(\mu) \mathbf{I}^{(l)}(-\mu, \phi, z_{l-1}) + \frac{1}{2\pi} \iint_{2\pi, \mu' > 0} \mathbf{R}^{\text{diff} \text{ top}, (l)}(\mu, \mu', \phi - \phi') \mathbf{I}^{(l)}(\mu', \phi', z_{l-1}) d\Omega' + \mathbf{I}^{\text{spec}, \text{bottom}, (l-1)}(\mu) \mathbf{I}^{(l-1)}(\mu, \phi, z_{l-1}) + \frac{1}{2\pi} \iint_{2\pi, \mu' < 0} \mathbf{I}^{\text{diff}, \text{bottom}, (l-1)}(\mu, \mu', \phi - \phi') \mathbf{I}^{(l-1)}(\mu', \phi', z_{l-1}) d\Omega'$$

**Computation:** 

- Step 1 a compute layer electromagnetic intrinsic proporties (Ke, Ks, Ka, P, eps) b - compute interfaces electromagnetic intrinsect proporties (R, T)
- Step 2 solve the radiative transfer equation

It's incredible how different they look:

	<b>Maximum extent</b> (aka traditional grain size), Dmax	<b>Correlation length</b> Exponential correlation fct A(x)	<b>Sphere radius</b> (distribution), stickiness: a, τ	<b>Sphere radius</b> (distribution), stickiness: a, τ
<u>Step1</u>	Empirical Ks Semi-empirical Ka	IBA W98 (Wahl=12) (Wahl<12)	DMRT	DMRT Short range Shih et al. 1997
<u>Step2</u>	кѕ,ка, q	кѕ,ка,Р(Ѳ)	кs,ка,Р(Ө	кѕ,ка,Р(Ө)
	1-flux	6-flux	N-stream (spline)	N-stream (DISORT, Jin 1994
	HUT	MEMLS	DMRT-QMS	DMRT-ML
	Fortran / Matlab	Matlab / Fortran	Matlab	Fortran/Python
	FMI	C. Mätzler & co	L. Tsang & co	LGGE (now IGE)

# **Comparison of models**

# **Numerical comparisons** showed that none of the models is significantly/always better than the others.

Tedesco et al. 2006, « Intercomparison of Electromagnetic Models for Passive Microwave Remote Sensing of Snow »

Tian, B. « Quantifying inter-comparison of the microwave emission model of layered snowpacks (MEMLS) and the multilayer dense media radiative transfer theory (DMRT) in modeling snow microwave radiance (IGARSS) », 2010

L. Brucker et al. 2011, thesis and « Modeling time series of microwave brightness temperature at Dome C, Antarctica, using vertically resolved snow temperature and microstructure measurements »

Roy et al. 2013, « Brightness temperature simulations of the Canadian seasonal snowpack driven by measurements of snow specific surface area »

Kwon, Y, « Error Characterization of Coupled Land Surface-Radiative Transfer Models for Snow Microwave Radiance Assimilation », 2015

Roy, A., A. Royer, O. St-Jean-Rondeau, B. Montpetit, G. Picard, A. Mavrovic, N. Marchand, and A. Langlois, Microwave snow emission modeling uncertainties in boreal and subarctic environments, The Cryosphere 10, 623-638, doi:10.5194/tc-10-623-2016, 2016

Sandells, M., Essery, R., Rutter, N., Wake, L., Leppänen, L., and Lemmetyinen, J.: Microstructure representation of snow in coupled snowpack and microwave emission models, The Cryosphere, 11, 229-246, tc-11-229-2017, 2017

Royer A., A. Roy, B. Montpetit, O. Saint-Jean-Rondeau, G. Picard, L. Brucker, and A. Langlois, Comparison of commonly-used microwave radiative transfer models for snow remote sensing. Remote Sensing of Environment, 190, 247—259, doi:10.1016/j.rse.2016.12.020, 2017

Performing a fair comparison is challenging because of the many different components and the different « grain size » metrics (microstructure).

# **Comparison of models**

This talk:

# Are existing snow microwave emission models so different?



#### Reconcilate:

- → the different electromagnetic theories
- → the different micro-structure representation used by these models
- $\rightarrow$  the different **solutions** of the radiative transfer equation

Löwe and Picard (TC, 2015) and Pan et al. (2016)

#### Guess who?

Löwe and Picard (TC, 2015)

$$\begin{split} K^2 &= k^2 + \frac{f_v(k_s^2 - k^2)}{1 + \frac{k_s^2 - k^2}{3K^2}(1 - f_v)} \begin{cases} 1 + i \frac{2(k_s^2 - k^2)Ka^3}{9\left[1 + \frac{(k_s^2 - k^2)}{3K^2}(1 - f_v)\right]} \times \\ \left[1 + 4\pi n \int_0^\infty dr \ r^2 \left(g(r) - 1\right)\right] \end{cases} \end{split}$$

$$\tilde{\omega} = \frac{2}{9} \frac{a^3 f_v}{\kappa_e} \left| \frac{k_s^2 - k^2}{1 + \frac{k_s^2 - k^2}{3K^2} (1 - f_v)} \right|^2 \left[ 1 + 4\pi n \int_0^\infty dr \ r^2 \left( g(r) - 1 \right) \right]$$

$$\gamma^{\text{bi}}(\mathbf{\hat{o}},\mathbf{\hat{i}}) = \frac{k_{\text{eff}}^4}{4\pi V} |\mathbf{F}_f|^2 \sin^2 \chi$$
$$= \nu (1-\nu) (\epsilon_2 - \epsilon_1)^2 K^2 I \cdot k^4 \sin^2 \chi \qquad I = \frac{1}{\alpha} \int_0^\infty A(x) x \sin(\alpha x) dx$$

$$\boldsymbol{\epsilon}_{\text{eff}} = \frac{2\boldsymbol{\epsilon}_1 - \boldsymbol{\epsilon}_2 + 3\nu(\boldsymbol{\epsilon}_2 - \boldsymbol{\epsilon}_1) + \sqrt{(2\boldsymbol{\epsilon}_1 - \boldsymbol{\epsilon}_2 + 3\nu(\boldsymbol{\epsilon}_2 - \boldsymbol{\epsilon}_1))^2 + 8\boldsymbol{\epsilon}_1\boldsymbol{\epsilon}_2}}{4}$$

#### "IBA" and "DMRT QCA-CP" theories in MEMLS and DMRT\*

Löwe and Picard, 2015, in the low frequency limit (<=37 GHz for most snow) for spherical scatters

- Effective medium permittivity/wavenumber:

$$\begin{split} \varepsilon_{\text{eff},0}^{\text{IBA}} &= \frac{2\varepsilon_1 - \varepsilon_2 + 3\phi_2 \left(\varepsilon_2 - \varepsilon_1\right)}{4} \\ &+ \frac{\sqrt{\left(2\varepsilon_1 - \varepsilon_2 + 3\phi_2 \left(\varepsilon_2 - \varepsilon_1\right)\right)^2 + 8\varepsilon_1 \varepsilon_2}}{4} \\ \varepsilon_{\text{eff},0}^{\text{QCA-CP}} &= \frac{\varepsilon_1 - \frac{\left(\varepsilon_2 - \varepsilon_1\right)}{3} \left(1 - 4\phi_2\right)}{2} \\ &+ \frac{\sqrt{\left(\varepsilon_1 - \frac{\left(\varepsilon_2 - \varepsilon_1\right)}{3} \left(1 - 4\phi_2\right)\right)^2 + 4\varepsilon_1 \frac{\left(\varepsilon_2 - \varepsilon_1\right)}{3} \left(1 - \phi_2\right)}}{2} \end{split}$$



Absorption formulations are identicalScattering coefficients:

$$\begin{split} \kappa_{\rm s}^{\rm IBA} &= \frac{2}{9} k_0^4 a^3 \phi_2 \left| \frac{\left(\varepsilon_2 - \varepsilon_1\right) \left(2\varepsilon_{\rm eff,0}^{\rm IBA} + \varepsilon_1\right)}{\left(2\varepsilon_{\rm eff,0}^{\rm IBA} + \varepsilon_2\right)} \right|^2 S(0) \\ \kappa_{\rm s}^{\rm QCA-CP} &= \frac{2}{9} k_0^4 a^3 \phi_2 \left| \frac{3\varepsilon_{\rm eff,0}^{\rm QCA-CP} \left(\varepsilon_2 - \varepsilon_1\right)}{3\varepsilon_{\rm eff,0}^{\rm QCA-CP} + \left(\varepsilon_2 - \varepsilon_1\right) \left(1 - \phi_2\right)} \right|^2 S(0). \end{split}$$

S(0) = snow micro-structure

#### **DMRT-QCA**

#### IBA, (+Bi-continuous DMRT)



When IBA uses Sticky Hard Sphere like DMRT instead of exponential autocorrelation:



#### Conclusion:

The main difference between MEMLS and DMRT family is the microstructure

HUT has semi-empirical formulation of scattering/extinction coefficient Grain size  $\mathsf{d}_0$ 

$$\kappa_e = 0.0018 f^{2.8} d_0^{2.0}$$
 where d<sub>0</sub>= 1.5 (1-exp(-1.5 D<sub>max</sub>))

$$\begin{split} \kappa_{\rm s}^{\rm IBA} &= \frac{2}{9} k_0^4 a^3 \phi_2 \left| \frac{\left(\varepsilon_2 - \varepsilon_1\right) \left(2\varepsilon_{\rm eff,0}^{\rm IBA} + \varepsilon_1\right)}{\left(2\varepsilon_{\rm eff,0}^{\rm IBA} + \varepsilon_2\right)} \right|^2 S(0) \\ \kappa_{\rm s}^{\rm QCA-CP} &= \frac{2}{9} k_0^4 a^3 \phi_2 \left| \frac{3\varepsilon_{\rm eff,0}^{\rm QCA-CP} \left(\varepsilon_2 - \varepsilon_1\right)}{3\varepsilon_{\rm eff,0}^{\rm QCA-CP} + \left(\varepsilon_2 - \varepsilon_1\right) \left(1 - \phi_2\right)} \right|^2 S(0). \end{split}$$

Using micro-structure images  $\rightarrow$  a geometrically-based relationship between  $\,D_{max}$  and  $p_{ex}$ 

Pan, Durand and co-authors, 2016



HUT and IBA have very different scattering coefficients !

Surprising because HUT and MEMLS are known to have good performance...

HUT: snow is a strongly forward scattering: q=0.96



IBA and QCA-CP : snow scattering is almost isotropic (Rayleigh or moderate Mie)



Radiative transfer equation:

$$\mu \frac{\partial \mathbf{I}(\mu, \phi, z)}{\partial z} = -\kappa_{e}(\mu, \phi, z) \cdot \mathbf{I}(\mu, \phi, z) + \frac{1}{4\pi} \iint_{4\pi} \mathbf{P}(\mu, \phi; \mu', \phi', z) \cdot \mathbf{I}(\mu', \phi', z) d\Omega' + \kappa_{a}(\mu, \phi, z) T(z)$$

Different formulations of Ke and P may lead to <u>exactly</u> the same RT equation (and exactly the same solution)

#### e.g.

C. Mitrescu, , G.L. Stephens, On similarity and scaling of the radiative transfer equation, Journal of Quantitative Spectroscopy and Radiative Transfer 86, 4, 387–394, 2004

H.C. van de Hulst, Multiple light scattering, Academic Press, New York, 1980

Joseph, Wiscombe, Weiman. The Delta-Eddington Approximation for Radiative Flux Transfer. Journal of the Atmospheric Sciences, 1976, 33, 2452-2459.

# Similarity theory

Visible in the two-flux theory: single scattering albedo  $\omega$  (~Ks) and asymmetry factor g (~P):

$$\omega$$
, g  $\omega'$ , g'  $\omega' = \frac{(1-f)\omega}{1-f\omega}$   $g' = \frac{g-f}{1-f}$  for any f

M-delta approximation: choose f to reduce the forward peak (0 < f < g )  $\,$ 



# **Similarity theory**

Pan, Durand and co-authors, 2016



Broad agreement once the comparison takes into account the different phase function shape → HUT has similar behavior as MEMLS despite huge apparent differences

# Conclusion

Back to the introductive questions, my opiniated response:

#### - Is this diversity apparent or profound?

- overall apparent

- all models converge to the "right" snow behaviour and give reasonable results (not always for physically correct reasons)

#### - Why such a diversity?

- historical
- different focus/approach

#### - Is this diversity beneficial or counter-productive for the community?

- <u>it has been beneficial</u> until many users started to be spend more time performing numerical inter-comparisons (incl. myself) than really using models to develop useful algorithms for end-users.

#### - What about the dual mode? Good or bad?

- it's time (as of 2015) to merge both because of dual mode missions and insitu datasets

#### - Why a new model?

# Conclusion

#### - Why a new model?

We don't need a new model (yet) but we need:

a repository of microwave community knowledge = merge all RT models / theories in one code base, one framework



SMRT is plane-parallel multi-layer radiative transfer model

$$\mu \frac{\partial \mathbf{I}(\mu, \phi, z)}{\partial z} = -\mathbf{\kappa}_{e}(\mu, \phi, z) \cdot \mathbf{I}(\mu, \phi, z) + \frac{1}{4\pi} \iint_{4\pi} \mathbf{P}(\mu, \phi; \mu', \phi', z) \cdot \mathbf{I}(\mu', \phi', z) \, d\Omega' + \mathbf{\kappa}_{a}(\mu, \phi, z) \, T(z)$$

It works as every other such model:

- 1- Define the snowpack
- 2- Compute scattering, absorption and effective permittivity in ever layer
- 3- Solve the RT equation with given boundary conditions (active or passive mode)
- 4- Show the results



SMRT is highly structured modular model





How does SMRT compute scattering, absorption, phase function, effective permittivity (EM model)?



How does SMRT "see" snow?



# A few options for the permittivity formulations of materials (ice, water, brine, wetice, ...)



How does it work in practice?

SMRT is coded in Python and makes use a lot of Python goodies. Very explicit function and parameter naming  $\rightarrow$  user friendly

	<pre>from smrt import make_snowpack, make_model, sensor_list</pre>	
Inputs	<pre># prepare inputs thickness = [100] corr_length = [50e-6] temperature = [270] density = [320]</pre>	
Snowpack	<pre># create the snowpack snowpack = make_snowpack(thickness=thickness,</pre>	
	density=density, temperature=temperature, corr_length=corr_length)	
Sensor config	<pre># create the sensor radiometer = sensor_list.amsre('37V')</pre>	
Choose model	<pre># create the model m = make_model("iba_original", "dort")</pre>	
Run	<pre># run the model result = m.run(radiometer, snowpack)</pre>	
Outputs	<pre># outputs print(result.TbV(), result.TbH())</pre>	

Here, SMRT behaves like MEMLS

It is very easy to explore different medium configuration, different EM models, difference permittivity equations.



medium = snowpack + ice\_column